

**On the Role of Inertial Instability in Stratosphere Troposphere Exchange
Near Midlatitude Cyclones**

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1 *Abstract.* In simulations of midlatitude cyclones with the University of Wisconsin
2 Nonhydrostatic Modeling System (UWNMS), mesoscale regions with large negative absolute
3 vorticity commonly occur in the upper troposphere and lower stratosphere (UTLS), overlying
4 thin layers of air with stratospheric values of ozone and potential vorticity (PV). These locally
5 enhanced stratosphere – troposphere exchange (STE) events are related to upstream convection
6 by tracing negative equivalent potential vorticity (EPV) anomalies along back trajectories. The
7 mutual coincidence of negative absolute vorticity, PV, and EPV, indicating inertial instability in
8 the UTLS, is shown to occur in association with enhanced STE signatures. Here results are
9 presented for two midlatitude cyclones in the upper Midwest where convection develops
10 between the subpolar and subtropical jets.

11 Mesoscale regions of negative EPV air originate upstream in the boundary layer. As they
12 are transported through convection, EPV becomes increasingly negative toward the tropopause.
13 In association with the arrival of each large negative EPV anomaly, a locally-enhanced poleward
14 surge of the subpolar jet occurs, characterized by high turbulent kinetic energy and low
15 Richardson number. Isosurfaces of wind speed show that gravity waves emanating from
16 inertially unstable regions modulate both jets simultaneously. It is shown that inertially unstable
17 convective outflow surges can facilitate STE locally by fostering poleward acceleration in the
18 UTLS, turbulent entrainment, and enhanced folding of tropospheric air over stratospheric air
19 underneath the subpolar jet.

20

21 **1. Introduction**

22 *a. Motivation*

23 Stratosphere – troposphere exchange (STE) of air in the upper troposphere and lower
24 stratosphere (UTLS) is an essential part of the Brewer-Dobson circulation, which circulates mass
25 between the stratosphere and troposphere, and therefore exerts a primary influence on the
26 distribution of climatically important constituents such as ozone, water vapor, and volcanic
27 aerosol. An abrupt increase in static stability occurs at the tropopause, caused by stratospheric
28 ozone heating, together with surface heating and buoyant adjustment by convection and
29 baroclinic waves in the troposphere (e.g., Manabe and Wetherald, 1967). The extratropical
30 tropopause is defined to occur when the lapse rate is less than 2 K/km, but is often defined by
31 values of Ertel’s potential vorticity ($PV = \frac{1}{\rho} \frac{\partial \theta}{\partial z} (f + \zeta)$, where ρ is density, θ is potential
32 temperature, f is the Coriolis parameter, and ζ is relative vorticity) in the range $\sim 1 - 4$ PVU ($1 \times$
33 $10^{-6} \text{ m}^2 \text{ K kg}^{-1} \text{ s}^{-1}$), or of ozone mixing ratio in the range 100-200 ppbv. According to the non-
34 transport theorem (Andrews et al., 1987), diabatic and frictional processes are required for
35 molecules to cross the tropopause (cf. Haynes et al., 1990). STE occurs in the vicinity of
36 thunderstorm tops, breaking Rossby waves, and breaking gravity waves, which cause mixing in
37 the UTLS (Holton et al., 1995). Estimates of STE vary by more than a factor of 2 (e.g., Stohl et
38 al., 2003), highlighting the need for improved understanding of the underlying dynamics.

39 The present work is motivated by our diagnosis of enhanced STE in conjunction with
40 regions of inertial instability in simulations of many different midlatitude cyclones with the
41 University of Wisconsin Nonhydrostatic Modeling System (UWNMS). These simulations were
42 carried out to investigate wintertime quasi-stationary precipitation features and to better
43 understand the origin of the thin sheets (\sim hundreds of meters thick) of stratospheric air, which
44 appear to curl poleward and downward into the troposphere near midlatitude jets in association
45 with midlatitude Rossby wave breaking (RWB) (e.g., Danielson, 1968). High resolution

46 dropwinsonde data sampling a wide variety of geographical locations and jet configurations,
47 including cutoff lows (Fig. 4, Shapiro, 1974), southeastward jets (Fig. 5, Shapiro, 1978; Fig. 3,
48 Shapiro, 1980), a digging trough (Fig. 1, Shapiro, 1981), an Icelandic low (Fig. 4, Shapiro, 1985)
49 and a staircase of westerly jets (Fig. 16, Shapiro et al., 1987), illustrate the ubiquitous sharpness
50 and typical structure of stratospheric intrusions. Aircraft campaigns, satellite data analysis, and
51 model analysis of trace constituents have shown that these sheets of air tend to be turbulent and
52 exhibit a mixed stratospheric / tropospheric chemical signal (e.g., Browell et al., 1989; Zahn et
53 al., 2000; Hoor et al., 2002; Stohl et al., 2003; Bowman et al., 2007; Pan et al., 2007; Hegglin
54 et al., 2008). The mystery of how thin stratospheric intrusions are formed strongly motivates
55 the present work.

56 In examining simulations of quasi-stationary, banded precipitation events in the upper
57 Midwest with the UWNMS, regions of negative equivalent potential vorticity (EPV) were found
58 near the tropopause on the anticyclonic shear side of the jet in the model. As will be
59 demonstrated, negative EPV in the UTLS (where it is statically stable and very dry) implies
60 inertial instability. Adjacent to these negative EPV regions were found layers of turbulently
61 mixed air associated with folds in PV and ozone. In this paper we explore the hypothesis that
62 mesoscale regions of inertial instability near the jet can facilitate STE. The proposed
63 mechanism for STE is that inertial instability causes poleward acceleration of uppermost
64 tropospheric air, leading to a locally enhanced fold, with increased turbulent entrainment of
65 stratospheric air.

66

67 *b. RWB, inertial instability, and westerly jets*

68 The subtropical westerly jets, which result from a buoyancy-driven redistribution of
69 angular momentum by local Hadley circulations, are collocated with strong baroclinic energy
70 conversion processes, as synoptic scale Rossby waves amplify, break, and are absorbed in the
71 UTLS (Andrews et al., 1987). Often a subpolar jet near 300 hPa (315-330 K) is present
72 simultaneously with a subtropical westerly jet near 200 hPa (340-350 K) (e.g., Christenson and
73 Martin, 2014). RWB occurs in 3D, with occlusion leading to a reversal of the normally positive
74 meridional gradient of PV, and tropopause folds. RWB implies a cascade of energy to smaller
75 scales, with small-scale mixing and radiative damping required for wave absorption (e.g.,
76 McIntyre and Palmer, 1983). Convective, Kelvin-Helmholtz, inertial, and baroclinic instabilities
77 tend to generate gravity waves. Gravity waves also amplify, break, and are absorbed through
78 small-scale mixing and radiative damping. Turbulence from breaking waves and instability is
79 required for STE.

80 A discussion of cyclonic and anticyclonic RWB, together with transport pathways in
81 midlatitude cyclones, is given by Thorncroft et al. (1993). The warm conveyor belt (WCB) can
82 split, with one branch curving cyclonically poleward, often riding up and over part of the lower
83 stratosphere, and plays a central role in the present work. This can contribute to the
84 identification of multiple tropopauses (e.g., Randel et al., 2007). The dry stratospheric conveyor
85 belt (DCB), which typically curves cyclonically equatorward into the troposphere, may aid the
86 formation of stratospheric intrusions by differential advection. Stratospheric intrusions often
87 take the form of elongated PV streamers as seen at constant height and thin sheets as seen in
88 cross-sections. These dry intrusions are sometimes related to severe wind events at the surface
89 (Browning and Reynolds, 1994) and forest fire outbreaks (Zimet et al., 2007; Schoeffler, 2013).

90 In quasi-geostrophic theory, a preferred region of poleward and downward circulation is
91 expected to occur near the jet entrance region (e.g., Markowski and Richardson, 2010). Cold air
92 advection along the UTLS jet axis promotes a preference for sinking downstream of a trough at
93 synoptic scales, via the Sutcliffe development mechanism (Martin, 2006; Lang and Martin,
94 2010). Sawyer (1949) explored the possible role of inertial or symmetric instability in the
95 UTLS in modifying midlatitude cyclogenesis, jet stream behavior, and meridional circulations.
96 It is now understood that inertial instability plays a role in establishing quasi-stationary
97 precipitation bands in midlatitude cyclones (e.g., Bennetts and Hoskins 1979; Knox 2003), and
98 phenomena on the edges of anticyclones (Stevens and Ciesielski 1986; Knox 1997). Jones and
99 Thorpe (1992) studied 3D inertial instabilities resulting from a region of negative PV and
100 concluded that the phenomenon is fundamentally the same as classical 2D symmetric instability.

101 Sato and Dunkerton (2002) showed that conditions are inertially unstable more than 20%
102 of the time in the upper troposphere on the equatorward side of the subtropical westerly jet south
103 of Japan during boreal winter. They found that the existence of distinctive layered perturbations
104 in the UTLS corresponded with times of inertial instability. Schumacher and Schultz (2000)
105 also found that inertially unstable conditions are common in the subtropical troposphere. Knox
106 and Harvey (2005) compiled a climatology of inertial instability and RWB and found that it is
107 inertially unstable over 2-5% of the midlatitude UTLS at any given time.

108 Planetary Rossby waves refracting into the tropical middle atmosphere encounter
109 inertially instable conditions, thereby exciting quasi-stationary 3D circulations, with vertical
110 motions leading to “pancake structures” in the temperature field (Hitchman et al. 1987; Hayashi
111 et al., 2002). As RWB proceeds, air with anomalous PV is advected across the equator, with
112 resultant inertial instability and overturning circulations. The associated divergence /

113 convergence fields ensure that the horizontal and vertical scales of inertial instability are smaller
114 than Rossby waves, facilitating an enstrophy cascade during RWB and helping to homogenize
115 PV (O’Sullivan and Hitchman, 1992).

116

117 *c. Convective influences on STE near jets*

118 A growing body of literature explores the role of latent heating variations during ascent in
119 the WCB in creating mid-tropospheric positive PV anomalies and upper tropospheric negative
120 PV anomalies (Pomroy and Thorpe, 2000; Knippertz and Wernli, 2010; Lang, 2011; Schemm et
121 al., 2013; Madonna et al., 2013). These papers guide our interpretation of heating along
122 trajectories through convection embedded within the WCB.

123 Cooper et al. (2004) showed that penetration of stratospheric intrusions by convection
124 facilitates subsequent mixing into the troposphere. Homeyer et al. (2011) found that convection
125 into stratospheric intrusions can yield a distinctive mixture of more than 125 ppmv H₂O and 100
126 ppbv O₃. Griffiths et al. (2000) argued that the PV anomaly associated with a fold can help to
127 induce convection. Conversely, Lang (2011) showed that convection can lead to intensification
128 of lower stratospheric fronts.

129 An example of ozone STE associated with a convective complex over China in March
130 2001 is described by Hitchman et al. (2004) and Kittaka et al. (2004), where STE occurred via
131 self-induced peripheral descent, similar to the circulation around an atomic explosion or
132 idealized warm bubble simulation (e.g., Fig. 5, Wicker and Skamarock, 1998). The convective
133 complex occurred in the westerly jet south of Japan, and thereby had access to high ozone in the
134 UTLS on the poleward side of the jet. Aircraft and UWNMS diagnoses showed that air with
135 high PV, high ozone, and very low water vapor curled downward and inward to ~6.5 km and

136 promoted the decay phase of the complex. It will be shown that, in the cases studied here,
137 enhanced stratospheric intrusions immediately follow convective injection of negative PV,
138 inertially unstable air at jet level.

139 Hoggatt and Knox (1998) studied an elevated convection event in the upper Midwest
140 during 14 July 1995 and discovered that elongated bands of light rain tend to coincide with a
141 narrow region of negative PV. Simulations of the event indicated a quadrupole checkerboard
142 pattern in the vertical motion field, signifying inertial instability, as in the equatorial middle
143 atmosphere. Schumacher et al. (2010) simulated quasi-stationary precipitation bands and found
144 that they occurred in an environment with a well-mixed baroclinic boundary layer, positive
145 convective available potential energy, and widespread negative PV. They argued that ascent
146 caused by frontogenesis and banded moist convection produced narrow regions of negative
147 absolute vorticity directly by the upward transport of low-momentum air. Convective bands
148 initiated within the ascending branch of the secondary circulation were associated with elevated
149 and near-surface frontogenesis. Their work is explicit about the presence of negative PV, and
150 inertial instability on the poleward side of the surface cyclone and anticyclonic side of the jet.
151 Their focus, however, was on the cause of the precipitation maximum and to compare simulated
152 versus observed convective bands near complex terrain.

153

154 *d. Organization of paper*

155 Dynamical diagnostic quantities are introduced in section 2, while the data and
156 simulations are described in section 3. In sections 4 and 5 two case studies with the UWNMS
157 are presented which feature pronounced, quasi-stationary, elongated precipitation bands common
158 in the upper Midwest during winter and spring. The 3D relationship among the locations of the

159 jets, negative EPV, meridional flow, upstream convection, and synoptic storm features is
 160 highlighted. The intimate temporal linkage between negative EPV anomalies, the signature of
 161 sharpened stratospheric intrusions, and gravity waves connecting the jets is shown in
 162 supplementary film loops. Case 2 features the signature of STE as seen in ozone initialized in
 163 the UWNMS with globally assimilated ozone from the Goddard Earth Observing System (GEOS
 164 ozone) (Stajner et al., 2008). A summary statement of the proposed mechanism, together with a
 165 schematic diagram, are given in section 6. It will be argued that inertial instability facilitates
 166 STE by causing “overfolds” in the UTLS, as part of the baroclinic energy conversion process
 167 arising from air masses of different density subject to gravity and rotation.

168

169 **2. Dynamical diagnostics**

170 *a. Inertial instability, absolute vorticity, and PV*

171 This section relates inertial instability to negative EPV and introduces the PV tendency
 172 equation. Inertial instability occurs when angular momentum decreases radially outward, such as
 173 in flow between two rotating cylinders (Rayleigh, 1916; Taylor, 1923). When this criterion for
 174 instability is met, parcels accelerate and rearrange themselves, tending to stabilize the angular
 175 momentum profile. Conversely, one might expect resistance to radial motion when angular
 176 momentum increases with radius. Knox (2003) provides a clear interpretation and
 177 comprehensive overview of inertial instability theory and related phenomena.

178 Consideration of radial parcel displacement δs in an axisymmetric vortex yields the
 179 oscillator equation

$$180 \quad \frac{\partial^2}{\partial t^2} \delta s = - f (f + \zeta) \delta s , \quad (1a)$$

181 with acceleration occurring (inertial instability) if the coefficient is negative:

182 $f (f + \zeta) < 0 ,$ (1b)

183 where $f = 2 \Omega \sin \varphi$ and $\zeta = \partial v / \partial x - \partial u / \partial y$ is relative vorticity (Eliassen and Kleinschmidt, 1957;
 184 Holton, 2006). From the reference from of the earth's rotation axis, absolute angular
 185 momentum per unit mass is given by $m = r u + r^2 \Omega$, where $r = a \cos \varphi$, a is earth radius, u is
 186 zonal flow, and Ω is the angular frequency of the earth's rotation. Inviscid inertial instability
 187 will occur in sufficiently strong anticyclonic relative vorticity, which can occur on the
 188 equatorward flank of a westerly jet. Inertial instability implies a divergence / convergence
 189 pattern associated with parcel acceleration and deceleration at the edges of inertially unstable
 190 regions, which in turn implies 3D circulations, enhanced turbulence, and mixing (O'Sullivan and
 191 Hitchman, 1992).

192 The inertial instability criterion is related to Ertel's potential vorticity (PV) and the
 193 gradient of angular momentum on an isentropic surface as follows:

194 $f \text{ PV} = \frac{1}{\rho} \frac{\partial \theta}{\partial z} f (f + \zeta) = - \frac{1}{r} \nabla_{\theta} m < 0 ,$ (2)

195 (e.g., Hoskins et al., 1985; Hitchman and Leovy, 1986). Inertial instability occurs when there is
 196 anomalously-signed PV for a given hemisphere, or anomalous absolute vorticity (first relation in
 197 (2)). It occurs when angular momentum increases toward the rotation axis (second relation in
 198 (2)). This criterion for instability is easier to satisfy near the equator where f is small and angular
 199 momentum surfaces are parallel to the surface of the earth. In the Northern Hemisphere, where f
 200 > 0 , the flow is inertially unstable if the absolute vorticity of the basic flow is negative: $f + \zeta < 0$.
 201 Note also that large values of PV (e.g., stratospheric) imply a stronger angular momentum
 202 gradient, hence resistance to poleward displacement.

203

204 *b. Equivalent potential vorticity*

205 Let us examine the implications of negative equivalent potential vorticity (EPV) in the
206 Northern Hemisphere:

$$207 \quad \text{EPV} = \frac{1}{\rho} \frac{\partial \theta_e}{\partial z} (f + \zeta) < 0 \quad , \quad (3a)$$

$$208 \quad \text{where} \quad \theta_e = \theta \exp\left(\frac{Lw_s}{c_p T}\right) \quad , \quad (3b)$$

209 $L = 2.5 \times 10^6 \text{ J kg}^{-1}$ is the latent heat of condensation, w_s is saturation water vapor mixing ratio,
210 $C_p = 1005 \text{ J kg}^{-1} \text{ K}^{-1}$ is specific heat at constant pressure, and T is temperature. θ_e allows for
211 energy conservation between vapor phase change and internal energy. According to (3b), in a
212 moist WCB starting at $\sim 12 \text{ g/kg}$ in the boundary layer, in the limit of complete elimination of
213 water vapor (a reduction of w_s to 0), θ would increase by $\sim 32 \text{ K}$ (e.g., Holton, 2006). In the
214 UTLS water vapor mixing ratios ($\sim 10^{-6}$) are reduced by four orders of magnitude relative to a
215 humid boundary layer ($\sim 10^{-2}$), implying that in the UTLS θ_e is very nearly equal to θ , and that
216 EPV is very nearly equal to PV. Moisture modifies static stability, especially in the boundary
217 layer, therefore the criterion $\partial \theta_e / \partial z < 0$ is the most useful for evaluating static instability.

218 Eqn. (3) shows that EPV is negative if the atmosphere is either moist statically unstable
219 or if it is inertially unstable. Air that is moist statically unstable that is brought by convection to
220 the base of the stratosphere will be statically stable but it may still have negative EPV, implying
221 that upstream convective instability can lead to downstream inertial instability.

222 In order to understand the origin of the negative EPV anomalies we consider the PV
223 tendency equation. Due to turbulent mixing and diabatic processes, PV will not be conserved
224 (Andrews et al., 1987). In warm conveyor belts and thunderstorms latent heating dominates.

225 Following Pomroy and Thorpe (2000), considering the dominant vertical component, and
226 ignoring frictional effects, the PV tendency equation becomes

$$227 \quad \frac{d PV}{dt} = \frac{1}{\rho} \left(\zeta \frac{\partial}{\partial z} \left(\frac{d\theta}{dt} \right) \right) . \quad (4)$$

228 In the cases to be shown, air found on the equatorward side of the subpolar westerly jet often
229 originates in the moist boundary layer of a subtropical air mass. In considering how this air
230 might be able to facilitate STE, it is important to include water vapor phase changes in defining
231 potential temperature used in calculating PV. Since convective updrafts occur where
232 $\partial\theta_e/\partial z < 0$, it is possible for air with negative EPV to be created and transported into the upper
233 troposphere, as dehydration causes EPV to asymptote to PV, preserving its negative value, with
234 static stability at the base of the stratosphere implying inertial instability. The upward decrease
235 of latent heating in the upper portion of thunderstorms can cause significant negative EPV
236 anomalies.

237 EPV dynamics are a vital contributor in assessing conditions conducive to conditional
238 symmetric instability (CSI). The concept of EPV has been widely used in studies of CSI in
239 baroclinic systems since CSI was first introduced as a potential mechanism for the generation of
240 frontal rainbands (Bennetts and Hoskins, 1979; Emanuel, 1979; Montgomery and Farrell, 1991;
241 Schumacher and Schultz, 2000; Schultz and Knox, 2007). EPV has proven useful in diagnosing
242 squall lines (Zhang and Cho, 1991), extratropical cyclones (Cao and Cho, 1995; Cao and Zhang,
243 2004; Brennan et al., 2007), convective snowstorms (Halcomb and Market, 2003), and
244 enhancing our understanding of the dynamics of the UTLS (Hoskins et al., 1985; Morgan, 1998).
245 The present paper focuses on implications of inertial instability for enhancing STE in the UTLS,
246 as highlighted by regions of large negative EPV.

247

248 *c. Turbulent kinetic energy and Richardson number*

249 Bernard (2013) studied geostrophic turbulence near rapid changes in stratification and
250 found that the forward cascade of buoyancy variance implies an enhanced region of turbulence
251 near the tropopause. Turbulence and gravity waves are noticeably enhanced in the UTLS near
252 westerly jets, and play an important role in mixing (Lindzen and Tung, 1976; Uccellini et al.,
253 1987; Pavelin et al., 2002; Whiteway et al., 2003; Duck and Whiteway, 2005; Koch et al.,
254 2005). Mixing near the jet stream at the tropopause and internal waves breaking in the vicinity
255 of the tropopause are potentially important in the exchange of trace gases such as ozone between
256 the troposphere and the stratosphere (Clayson and Kantha, 2007).

257 The TKE and Ri patterns to be shown highlight regions of small-scale mixing, which
258 help to exchange air between the stratosphere and troposphere. Kinetic energy per unit mass
259 (TKE, $\varepsilon \sim \text{m}^2 \cdot \text{s}^{-2}$) is a measure of the intensity of turbulence and is defined as

260
$$\varepsilon = \frac{1}{2} (\overline{u'^2} + \overline{v'^2} + \overline{w'^2}), \quad (5)$$

261 where the overbar indicates time and space averaging of the turbulent eddy variances. The
262 UWNMS contains a prognostic equation for TKE, including advection, buoyancy generation,
263 shear generation, and turbulent dissipation (Tripoli, 1992a, b).

264 Distributions of the Richardson number, Ri, highlight where convective or shear
265 instability is generating TKE. Ri is the ratio of the mechanical generation of TKE by wind shear
266 to the buoyancy production or inhibition by static stability:

267
$$\text{Ri} = \frac{\frac{g}{\theta} \frac{\partial \theta}{\partial z}}{\left(\frac{\partial u}{\partial z}\right)^2 + \left(\frac{\partial v}{\partial z}\right)^2}, \quad (6)$$

268 where the numerator is equal to the square of the buoyancy frequency. The flow is dynamically
269 unstable for flows with $Ri < 0.25$ (e.g., Markowski and Richardson, 2010). Once generated,
270 turbulent flow tends to stay turbulent, even for Ri numbers as large as 1.0 (e.g., Stull, 1988).

271

272 **3. Data and Analysis Methods**

273 During the course of simulating banded precipitation maxima in midlatitude cyclones
274 with the UWNMS, distinctive mesoscale regions of negative EPV were often observed in the
275 UTLS on the equatorward side of the subpolar jet. The spatial and temporal coincidence of
276 pronounced mesoscale stratospheric intrusions on the poleward side of the jet, as seen in PV, θ ,
277 and ozone, suggested that perhaps the regions of negative PV were related to the locally
278 enhanced STE. During 6-8 February 2008 (Case 1), phenomena of interest were located over
279 south-central Wisconsin and northern Illinois. During 22-23 April, 2005 (Case 2), negative EPV
280 and local overturning circulations occurred over the Ohio Valley.

281 Each case is modeled with the UWNMS (Tripoli 1992a, b), initialized with 2.5°
282 European Centre for Medium-range Weather Forecasts (ECMWF) data. Studies of processes in
283 the UTLS, including STE, using the UWNMS include Pokrandt et al. (1996), Hitchman et al.
284 (1999; 2001; 2004), Kittaka et al. (2004), and Bükér et al. (2005; 2008). The specified UWNMS
285 resolution for each case is 20 km x 20 km x 300 m, with a grid volume of 152 x 152 x 60 points.
286 The top of the model is set to 16 km, with a 1500 m Rayleigh sponge layer. Each of the
287 simulations was initialized ~24 hours before the phenomena of interest began and continued for
288 48 hours.

289 Case 2 was chosen in part because of the availability of high quality globally assimilated
290 ozone fields from GEOS data for 2005 (Stajner, 2008). GEOS ozone data are used to initialize

291 the UWNMS, and to update ozone on the UWNMS side boundaries. Detailed ozone structures
292 evolve in the interior of the UWNMS over the subsequent 48 hours.

293 Vis5D is used to view the 3D structure of each midlatitude cyclone and the associated
294 dynamical quantities in the UWNMS, and to create trajectories which connect the negative EPV
295 anomalies with upstream convection. National Centers for Environmental Prediction (NCEP)
296 Eta model analyses with 80 km resolution are used to illustrate the synoptic context of the two
297 cases. They provide corroborating evidence of the prevalence of inertial instability on the
298 equatorward side of the subpolar jet.

299

300 **4. Case Study 1 (6-8 February 2008)**

301 *a. Synoptic overview*

302 This destructive mid-winter storm during 6-8 February 2008 featured quasi-stationary
303 snowbands, which trapped over 2000 vehicles for 24 hours on the I-90 freeway in southern
304 Wisconsin, and spawned deadly tornadoes along the cold front throughout the Mississippi and
305 Ohio Valleys. In addition to chronic negative absolute vorticity on the equatorward side of the
306 subpolar jet, convection erupted along the cold front, injecting air into the UTLS between the
307 subtropical and subpolar jets, generating EPV values of less than -3 PVU on the equatorward
308 flank of the subpolar jet.

309 Figure 1 shows an overview of the synoptic scale features in the lower troposphere
310 associated with the storm at 0000 UT 6 February 2008, using NCEP Eta model reanalyses.
311 Figures 1a and 1b show the surface low pressure centered over southern Missouri and the
312 associated baroclinic zone at 850 hPa. During the low's migration northeastward over Ohio,
313 quasi-stationary snowbands matured and persisted in southwestern Wisconsin and Northern

314 Illinois, with a characteristic east-west arc-shaped linear configuration, and snowfall in the range
315 12-18” (exceeding 1” water equivalent in Fig. 1c).

316 Figure 2 depicts the distribution of PV and wind speed at increasing levels in the UTLS
317 for the same time as in Fig. 1. At 300 hPa a strong subpolar jet with speeds in excess of 70 m/s
318 can be seen in Fig. 2f, arcing from Texas toward the Maritime Provinces, while at 200 and 150
319 hPa, (Figs. 2d, b), the subtropical jet can be seen extending from Eastern Mexico toward the
320 Mid-Atlantic states. A pronounced southward extension of high polar stratospheric PV air in the
321 300-250 hPa layer is shown in purple in Fig. 1e, along with a broad anticyclonic arc of negative
322 EPV in dark blue on the equatorward side of the subpolar jet (compare Figs. 2e, f). The pattern
323 of EPV in the 250-200 hPa layer includes negative regions over Illinois and Southern Quebec on
324 the equatorward side of the subpolar jet (compare Figs. 2c, d). The negative PV regions over
325 the Gulf of Mexico and Northern Florida occur on the equatorward side of the subtropical jet and
326 become more extensive in area upon ascending from the 300-250 hPa layer to the 200-150 hPa
327 layer (Figs. 2a, c, e).

328 A strong spatial correlation exists between the region of negative PV on the equatorward
329 side of the subtropical jet (Figs. 2e, f) and 24-hr accumulated precipitation during 0000 UT 6
330 February to 0000 UT on 7 February (Fig. 1c), consistent with previous work on symmetric
331 instability and regions of enhanced rainfall.

332

333 *b. Overview of UWNMS Case 1*

334 Figure 3 shows an oblique view from the southeast of a section from southeastern North
335 Dakota to western Virginia, in the UWNMS simulation at 2230 UT 6 February 2008. The
336 distribution of EPV in the section is shown in panels a, b, and d. Note the regions of negative

337 EPV in the upper troposphere over Wisconsin and Illinois, adjacent to stratospheric air in the
338 range 1-6 PVU over Minnesota. Figure 3c shows the same EPV contours but also includes a red
339 60 m/s speed isosurface and blue isosurfaces of negative EPV exceeding -2 PVU, excluding
340 values from behind the section. Note the coincidence of the subpolar jet and region of inertial
341 instability immediately adjacent to the stratospheric air. Note also the separation between the
342 subpolar jet over Iowa and the higher-altitude subtropical jet over Kentucky in this section.

343 Figure 3b shows the relationship between the EPV distribution, regions of EPV
344 exceeding -2 PVU on the equatorward side of the jet, as well as connections (back trajectory
345 ribbons) to the surface via the WCB and convection along the cold front (4 cm/s green updraft
346 isosurfaces). The regions of strongly negative EPV are seen to originate in the boundary layer
347 over Kansas ~14 hours earlier. A quantitative analysis of EPV evolution along the trajectory is
348 given in the next subsection. The poleward buckling of EPV contours in the UTLS and
349 enhanced signature of a stratospheric intrusion occurs with the arrival of inertially unstable air
350 from convective outflow in the WCB.

351 The contemporaneous distribution of TKE (blue contours) and Richardson number (black
352 contours) is shown in Fig. 3c. The high vault of negative EPV in the uppermost troposphere
353 generally coincides with low Ri and high TKE, with strong gradients along the base of the
354 stratosphere over the jet core and back into the troposphere along the upper edge of the
355 stratospheric intrusion. Note also the coincidence of TKE exceeding $1.0 \text{ m}^2 \text{ s}^{-2}$ and $\text{Ri} < 1$ in
356 the convective element found in the section over Illinois (Figs. 3b, c).

357 The evolution of PV and wind fields during the 10 hours preceding Fig. 3 is shown in
358 Fig. 4. At 1300 UT on 6 February 2008 a broad jet with winds exceeding 60 m/s occurred in the
359 UTLS over Iowa (Fig. 4a), with the tropopause sloping gently and monotonically upward toward

360 the south, as indicated by the 1 PVU contour (Fig. 4a) and transition to red (Fig. 4b). At this time
361 there were only modest regions of negative EPV near the jet (Fig. 4a). But by 1830 UT,
362 maximum wind speeds had increased to 75 m/s and shifted poleward, in conjunction with the
363 arrival of convective outflow air with PV < - 2 PVU (Figs. 4c, d). By 2100 UT the region of
364 very negative EPV had passed east of the section (Figs. 4e, f). This sequence shows that the
365 tropopause undergoes a distinct folding process, giving the appearance of a poleward and
366 downward circulation of stratospheric air into the troposphere underneath the subpolar jet.

367 The evolution of the jets, PV, angular momentum, and divergence during 1830 – 2230
368 UT on 6 February is shown in Fig. 5. At 1830 UT (panel a), high values of angular momentum
369 divided by earth's radius (black contours), are seen to bulge poleward in the UTLS, with the
370 maximum poleward displacement coinciding with the jet maximum (purple contours) and with a
371 region of strong inertial instability (EPV < -2 PVU). The region of actual inertial instability is
372 much larger, as evidenced by the area in which contours of m reverse sign, increasing poleward
373 in the UTLS from Missouri to Minnesota in this section.

374 The 60 m/s jet isosurface at 2000 UT is highlighted in light blue in Fig. 5b. New regions
375 of strongly negative EPV are approaching the section from the southeast, as old ones recede to
376 the east behind the jet isosurface. Note the large region where $\partial m / \partial y > 0$ in the upper
377 troposphere, tilting poleward with altitude into the jet. The structure of the poleward surge at
378 2130 UT is seen in Fig. 5c, with the greatest poleward motion occurring in the UTLS, coincident
379 with the greatest poleward excursion of tropical m contours. This poleward surge is a blend of
380 air from the subtropics and from convection along the cold front. This surge extends well into
381 the stratosphere, while an equatorward flow exists in the lower troposphere, with the boundary
382 between the two air masses tilting poleward with height. Differential advection in altitude can

383 readily produce the signature of a stratospheric fold. The poleward convective outflow surge
384 characterized by inertial instability overrides the layer of stratospheric air. This signature of
385 rising warm air and sinking cold air is to be expected for baroclinic synoptic waves. Note also
386 the enhanced region of strong inertial stability, where $-\partial m/\partial y > 0$, at the head of the poleward
387 surge in panels a-c.

388 The air at the top of the updraft (Fig. 5c) is strongly divergent in the UTLS (dashed
389 contours in Fig. 5d) and is coincident with a region of negative EPV. Theory predicts poleward
390 acceleration, hence horizontal divergence, in a region of inertial instability and initial poleward
391 motion. The striking pattern of divergence and convergence in the troposphere in Fig. 5d is
392 related to inertial instability and the generation of gravity waves, to be discussed below.

393 An animation of this process is provided in Supplemental movie S1, which highlights the
394 intimate temporal relationship between the arrival of negative EPV anomalies in the UTLS, the
395 poleward surges emanating from the regions of negative EPV, and enhanced signature of a
396 stratospheric intrusion resulting from overriding UTLS air. Figures 3-5 and S1 show that regions
397 of negative EPV within the jet originate in upstream convection and are associated with
398 poleward surges in the westerly jet. This association between negative absolute vorticity and
399 turbulence variables suggests that inertial instability is associated with enhanced mixing of air, as
400 well as with an enhanced intrusion signature. This provides evidence for the hypothesis that
401 strong vertical motion associated with convection creates negative EPV air in the uppermost
402 troposphere, with the resulting inertial instability favoring poleward flow and locally-enhanced
403 stratospheric intrusions.

404

405 *c. Changes of EPV along the path*

406 Negative anomalies of EPV can form due to diabatic processes associated with surface
407 cooling or with latent heat release in regions of ascent (Morgan, 1997). In convection embedded
408 in the warm upglide sector, a latent heating maximum can generate negative PV through vortex
409 shrinking (Joos and Wernli, 2012). Figures 2a-c show how updrafts transport unstable air from
410 the surface into the the jet, sometimes yielding regions of negative EPV exceeding -3 PVU at
411 heights near 11 km. The negative EPV air found equatorward of the jet initially resided near the
412 surface before being transported to the base of the stratosphere via updrafts exceeding 4 cm/s.
413 These negative EPV anomalies are typically ~1 km thick (3-4 vertical grid points), but the
414 horizontal scale varies considerably: ~40-200 km (2-10 grid points). Back trajectories in Fig. 3
415 reveal the origin of the negative EPV anomalies as they commenced in the unstable boundary
416 layer and became increasingly more negative while being lifted to the tropopause.

417 The probe function in Vis5d allows for sampling model variables along the trajectories in
418 this WCB region, including changes of EPV and net diabatic heating along the path.

419 A negative EPV anomaly typically originates near the surface with values around 0 PVU. As the
420 air rises into the upper troposphere, EPV becomes more negative, exceeding -3 PVU on the
421 equatorward side of the subpolar jet. By considering eqns. 3 and 4 one may estimate changes in
422 PV due to the pattern of latent heating. The density factor favors larger EPV anomalies in the
423 upper troposphere. Since latent heating maximizes in the mid-troposphere, then $\frac{\partial}{\partial z} \left(\frac{d\theta}{dt} \right) < 0$ in
424 the upper troposphere, generating negative PV. If a latent heating maximum of 20 K/ 6 hr is
425 assumed near 600 hPa, eqn. (4) suggests that with $\rho \sim 0.3 \text{ kg m}^{-3}$, $f \sim 10^{-4} \text{ s}^{-1}$, and $\delta z \sim 4000 \text{ m}$,
426 an EPV anomaly of $\sim -3 \text{ PVU}$ could be generated approaching the 200 hPa level. This
427 calculation supports the mechanism of generating negative EPV described by previous authors.

428

429 *d. Evolution and comparison of PV, EPV, and absolute vorticity at 8 km*

430 The horizontal signature of RWB in the upper troposphere may be seen in Fig. 4, which
431 quantitatively intercompares PV, EPV, and absolute vorticity at 8 km. The color at left is EPV,
432 while color at right is absolute vorticity, with stratospheric values shown in red and inertially
433 unstable regions in blue. Contours of PV (black) and EPV (white) are also shown in each panel.
434 The evolution of PV over a five hour period is shown, with the upper panels corresponding to
435 2230 UT 6 February, while the lower panels are for 0330 UT 7 February 2008. A close
436 comparison of the black PV contours and white EPV contours confirms that water vapor in the
437 UTLS makes only a minor quantitative difference between PV and EPV. This is consistent with
438 the reduction of water vapor concentration by four orders of magnitude in going from the surface
439 to the tropopause. The remarkable spatial coincidence of negative regions (blue) of EPV and
440 absolute vorticity shows that in the UTLS, each variable is a reliable indicator of regions of
441 inertial instability, hence may be useful in anticipating mesoscale motions such as radiation of
442 internal gravity waves.

443

444 *e. A case of “jet suturing” via gravity waves excited by inertial instability?*

445 Figures 7 and 8 show the evolution of the jet pair from the southwest and east,
446 respectively. Supplemental movies S2 and S3 show the evolution of this jet pair for the 48 hour
447 simulation as seen from the west and the east. In this sequence, poleward and upward moving
448 streams of air with $EPV < -1$ PVU are highlighted in aquamarine, while the 55 m/s speed
449 isosurface is shown in yellow. The right hand panels in Fig. 7 include horizontal sections of
450 absolute vorticity at 5 km. The air masses with large negative EPV are seen to travel upward
451 through an elongated band of negative absolute vorticity at 5 km, connecting to downward-

452 extending fingerlike structures of high wind speed from both jets. This remarkable mesoscale
453 disturbance is related directly to the convection. The fingers of high wind speed appear to have
454 their roots in the region of inertial instability. This mesoscale wave pattern fans out in the
455 vertical, as suggested by the convergence/divergence pattern in Fig. 6d, embracing both jets in
456 the UTLS.

457 One possible explanation is that the quasi-linear gravity wave-like field perpendicular to
458 the jets is due to inertia-gravity waves generated by the convection and inertial instability, which
459 have a preferred propagation direction antiparallel to the prevailing wind (Dunkerton and
460 Butchart, 1984). Eddy winds associated with gravity waves on the order of 10 m/s would be
461 sufficient to modulate the jet isosurface such as to cause this rib-like structure which connects
462 and modulates the two jets simultaneously, a stitching-together or suturing. This suggests that
463 gravity waves excited by convection and inertial instability may influence the co-evolution of
464 proximal subpolar and subtropical jets.

465 Figure 8 shows a view from the east of the ascending negative PV streamers merging
466 with downward-extending “stalactites” of high wind speed air. Note the poleward and
467 downward expansion of the jet isosurface from panel a through c. Since the jet core exits the
468 eastern model boundary, one may view snapshots of dollops of air with negative PV as they fly
469 poleward through the subpolar jet, with each event causing a poleward and downward extension
470 of the jet (cf. supplemental movie S3). A salient aspect of supplemental movies S2 and S3 is the
471 appearance of a suturing of the two jets. To the extent that gravity waves break and mix, this
472 may represent a mechanism of interaction between jets generated by convection and inertial
473 instability.

474

475 **5. Case Study 2 (22 April 2005)**

476 *a. Synoptic overview*

477 The spring storm of 22 April 2005 which caused significant precipitation across the Ohio
478 Valley also occurred between subpolar and subtropical jets, ahead of an equatorward-extending
479 stratospheric PV streamer. Even though wind speeds were weaker than in Case 1, EPV still
480 reached values of less than -2 PVU near 13 km. At 1200 UT on 22 April a moderate low
481 pressure system was found over the Ohio Valley, advancing toward the northeast (Fig. 9a).
482 Banded precipitation consistently formed and dissipated on the northwest side of the low
483 pressure center, lingering into the late evening hours of 22 April 2005 (Fig. 9b). A subtropical
484 westerly jet (Fig. 9c) is seen extending from Southern California over Texas and toward the Gulf
485 of Mexico, with notable regions of inertial instability (negative PV in Fig. 9d). Meanwhile, the
486 subpolar westerly jet extended from the Great Lakes toward Eastern Canada, with maximum
487 wind speeds near 50 m/s (Fig. 9d). During the time of the most intense precipitation, a distinct
488 patch of negative EPV (Figs. 9e, f) was found on the anticyclonic shear side of the subpolar jet,
489 consistent with the hypothesis that inertial instability enhances rainfall in midlatitude cyclones.
490 The existence of inertially unstable regions equatorward of westerly jets appears to occur
491 frequently, with convection between two jets particularly conducive for inertial instability-driven
492 poleward acceleration.

493

494 *b. UWNMS simulation of Case 2 with GEOS ozone*

495 The evolution of EPV (left) and GEOS ozone mixing ratio (right) is shown, together with
496 streamfunction in a meridional section extending from Georgia to southern Ontario, during 1200
497 - 2230 UT 22 April 2005 (Fig. 10). EPV isosurfaces of -2 PVU are shown in green. High

498 values of PV and ozone are indicated in red. Note the detailed agreement between the patterns of
499 stratospheric PV and stratospheric ozone. The mesoscale details of poleward deformation and
500 stratospheric intrusion are created by the model physics in the UWNMS from coarser initial
501 fields. Inertially unstable regions in the upper troposphere occur immediately above and
502 equatorward of stratospheric intrusions. The association of negative PV anomalies in the upper
503 troposphere and mesoscale disturbances in lower stratospheric PV is notable in panels c-f. The
504 narrow stratospheric intrusion over Indiana shows up clearly in both ozone and PV, coinciding
505 with the arrival of air with substantially negative PV at jet levels. As with Case 1, the
506 meridional streamfunction depicts a deep layer of poleward moving air in the UTLS, which
507 emanates from the convectively and inertially unstable updraft.

508 The convective updraft over Indiana (green streamfunction) exhibits the canonical
509 signature of positive PV production due to latent heating increasing with height, hence vortex
510 stretching in the lower troposphere, and production of negative PV in the upper troposphere due
511 to vortex compression (Figs. 10a, c, e). The same updraft also brought low values of ozone
512 upward into the UTLS (note the upward extension of light blue in Figs. 10b, d, f).

513 The development of sharp ozone structures in the UWNMS associated with negative PV
514 regions lends further support to the idea that inertial instability can facilitate stratospheric
515 intrusions. The structure in ozone and PV is consistent with upper tropospheric air surging
516 poleward over stratospheric air, yielding a trailing pattern of “stratospheric intrusion”
517 underneath.

518

519 **6. Summary of proposed mechanism**

520 Two case studies have been presented which provide evidence for a causal relationship
521 between inertial instability in the UTLS and locally-enhanced STE near midlatitude westerly
522 jets. This phenomenon occurs when a surface low pressure system, with associated convection
523 along a cold frontal boundary, exists between a subtropical and subpolar jet, with WCB updrafts
524 which lift the air to the base of the stratosphere, resulting in areas of enhanced inertial instability
525 near the subpolar jet. It was shown that negative regions of EPV, PV, and absolute vorticity
526 coincide in the UTLS, each serving equally well to indicate inertial instability. This inertial
527 instability facilitates poleward surges of air in the uppermost troposphere, yielding strong shears
528 at the base of the stratosphere, sufficient to give rise to a layer of low Ri and high TKE, which
529 bounds the layer of mixed air that surges poleward over the extratropical troposphere.

530 A mundane yet perhaps surprising result is that a “stratospheric intrusion” can be created
531 simply by a vertically-limited poleward surge in the UTLS, as shown schematically in Fig. 11.
532 There is no need to stipulate a circulation around a jet to achieve this structure. The poleward
533 expansion of air in the uppermost troposphere over the underlying stratospheric air creates a
534 stratospheric intrusion simply by kinematic folding. In these cases the poleward motion of
535 rising light air advects the subtropical westerly jet northeastward while the equatorward sinking
536 cold air in the lower troposphere tends to cause the subpolar jet to move southeastward. The
537 gravity wave field excited by convective / inertial instability between the two jets modulates the
538 two jets simultaneously, forming a rib-like connection between the two. We find that inertial
539 instability can act to accentuate poleward displacement in the WCB, promoting the large scale
540 baroclinic energy conversion by gravitational adjustment of air with differing density, subject to
541 the rigidity of rotation. The poleward extension of the subtropical westerly jet is resisted by the

542 enhanced meridional gradient of angular momentum (positive PV) of the extratropical lower
543 stratosphere.

544 Further work is required to diagnose the “suture” signature of jet merger associated with
545 inertial instability in the UTLS. We are currently exploring the seasonality and distribution of
546 inertial instability influences on STE. A companion paper emphasizes the relationship between
547 inertial instability, divergence, poleward surges, and mesoscale jet “flare-ups”.

548

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555

556

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Figure Captions

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781 Figure 3. The relationship among convection, inertial instability, jet locations, PV and
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784 subtropical jet over Kentucky. Panels a), b), and c) show black EPV contours (interval 1 PVU).
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786 isosurface, the blue (-2 PVU) isosurfaces indicate regions of strong inertial instability in the
787 UTLS, and the ribbons show 22.5-hr back-trajectories from these -2 PVU regions to where they
788 spent the first 10 hours in the unstable boundary layer over Kansas. Panel d) shows the two jets
789 (red 60 m/s isosurface) and the same blue -2 PVU isosurface, but with values behind the section
790 not shown. Panel c) shows the vault of low Ri (black contours, interval 0.5, range 0 - 3), regions
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795 black (white), interval 1 PVU. For each time, EPV is shown in color at the left, while absolute
796 vorticity (units 10^{-5} s^{-1}) is shown in color at the right. The color ranges are blue (< 0 ; inertially
797 unstable), yellow (0 - 1), and red (> 1 ; stratospheric air).

798 Figure 5. Meridional sections showing the relationship between absolute angular momentum,
799 wind speed, regions of inertial instability, divergence, and meridional streamfunction, as seen
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803 every 10 m/s), together with regions of highly inertially unstable regions (EPV < -2 PVU in
804 purple). Panel b) shows the co-location of the poleward-protruding high angular momentum air
805 mass and the 60 m/s wind speed isosurface. Panel c) shows meridional streamfunction in green,
806 indicating the depth of the poleward intrusion. Panel d) shows regions of convergence (solid
807 contours) and divergence (dashed contours), interval $10 \times 10^{-5} \text{ s}^{-1}$.

808 Figure 6. Onset of the effects of inertial instability as seen in vertical sections of PV, wind
809 speed, and regions of inertial instability in the UWNMS for 1300 UT (panels a and b), 1830 UT
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811 contours (interval 1 PVU), and in color in the right column (blue < 0 ; inertially unstable, yellow
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814 and d), and exiting behind the plane in panels e) and f).

815

816 Figure 7. View from the southwest of the influence of convection, inertial instability, and
817 gravity waves on the subtropical and subpolar jets (55 m/s speed isosurfaces in yellow) in the
818 UWNMS at 2230 UT 5 February (panels a and b), 0330 UT 6 February (panels c and d), and
819 0600 UT 6 February 2008 (panels e and f). The -1 PVU isosurfaces are shown in light green in
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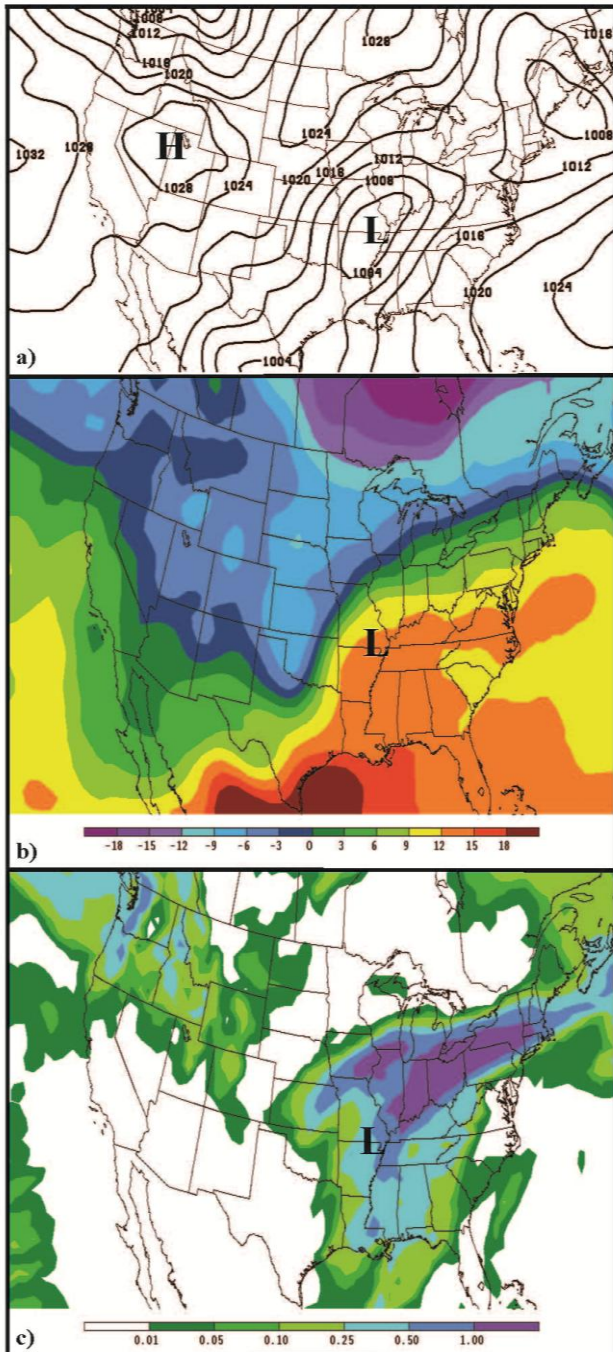
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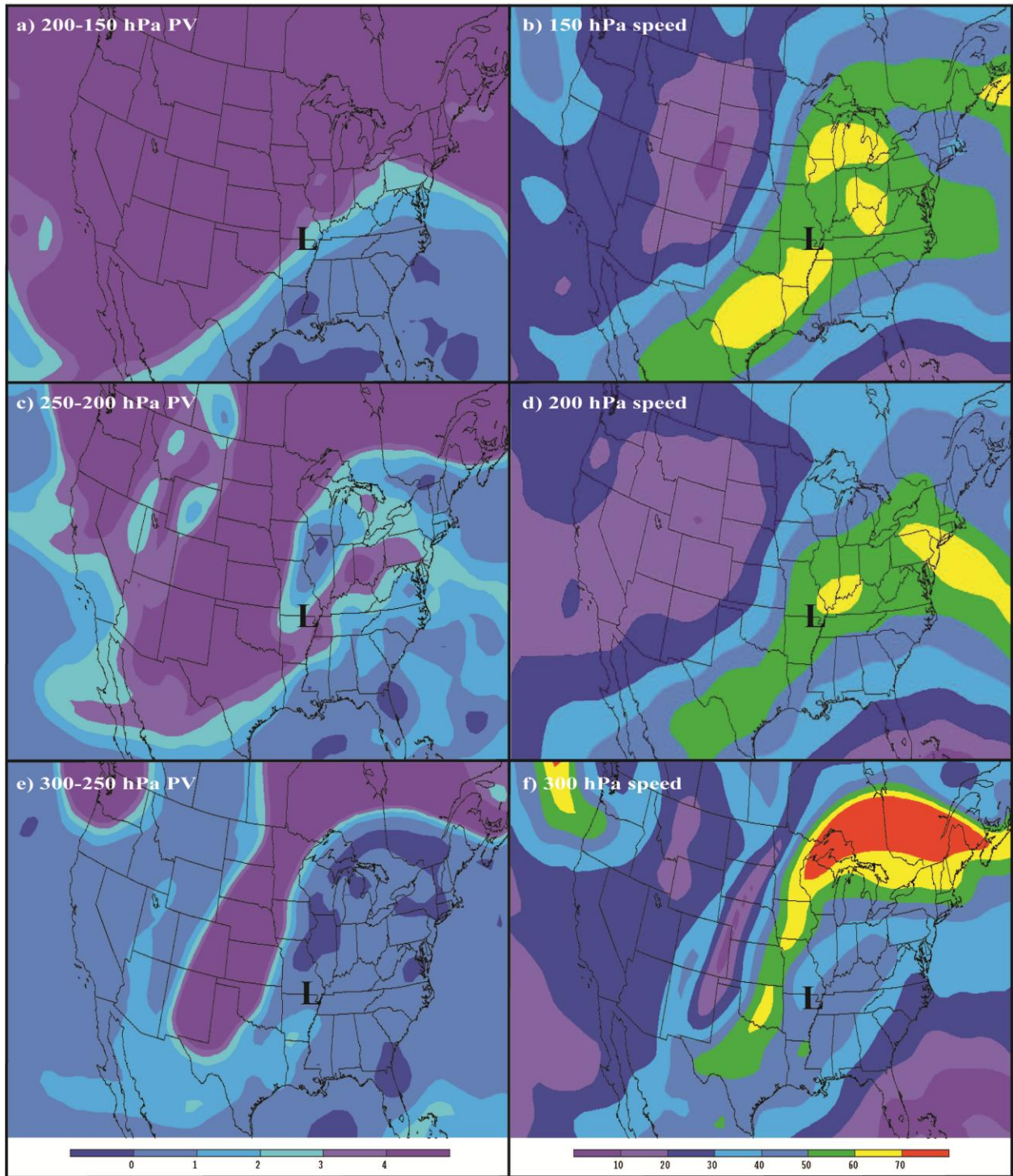
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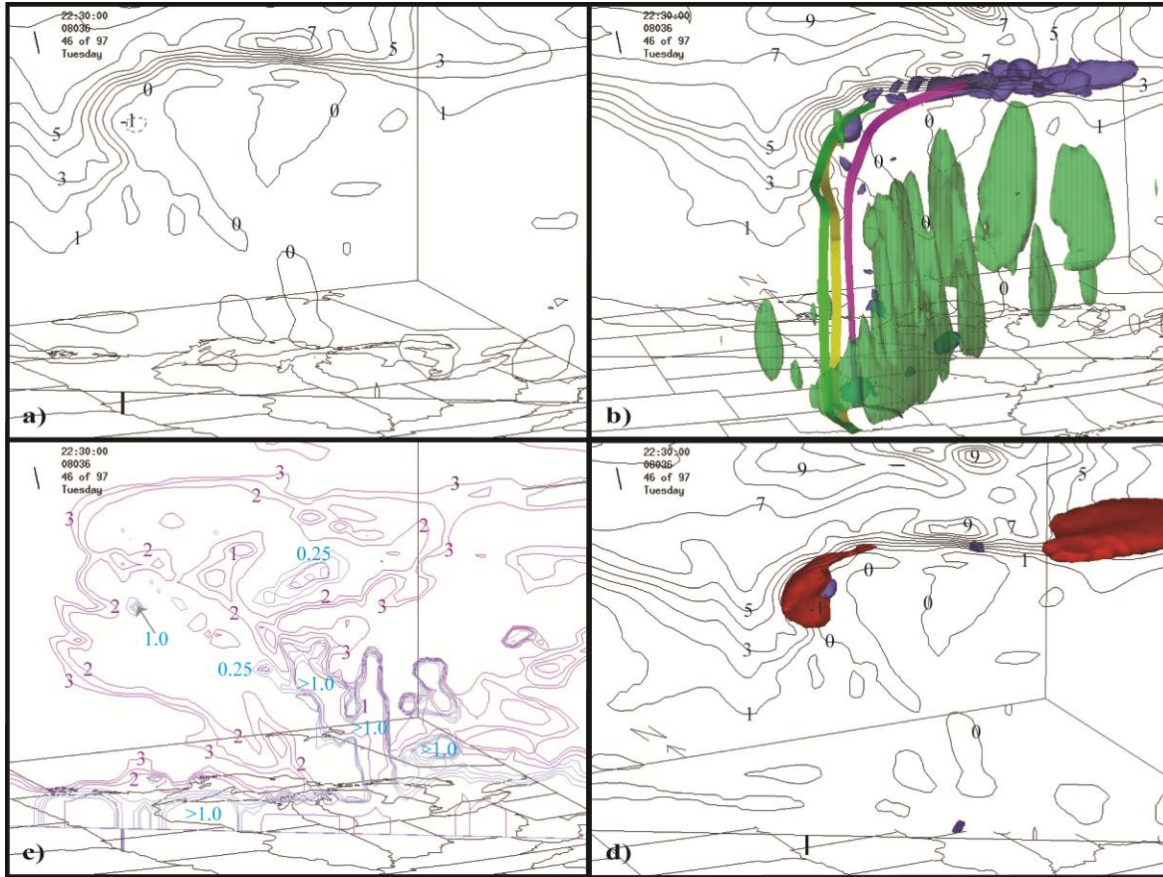
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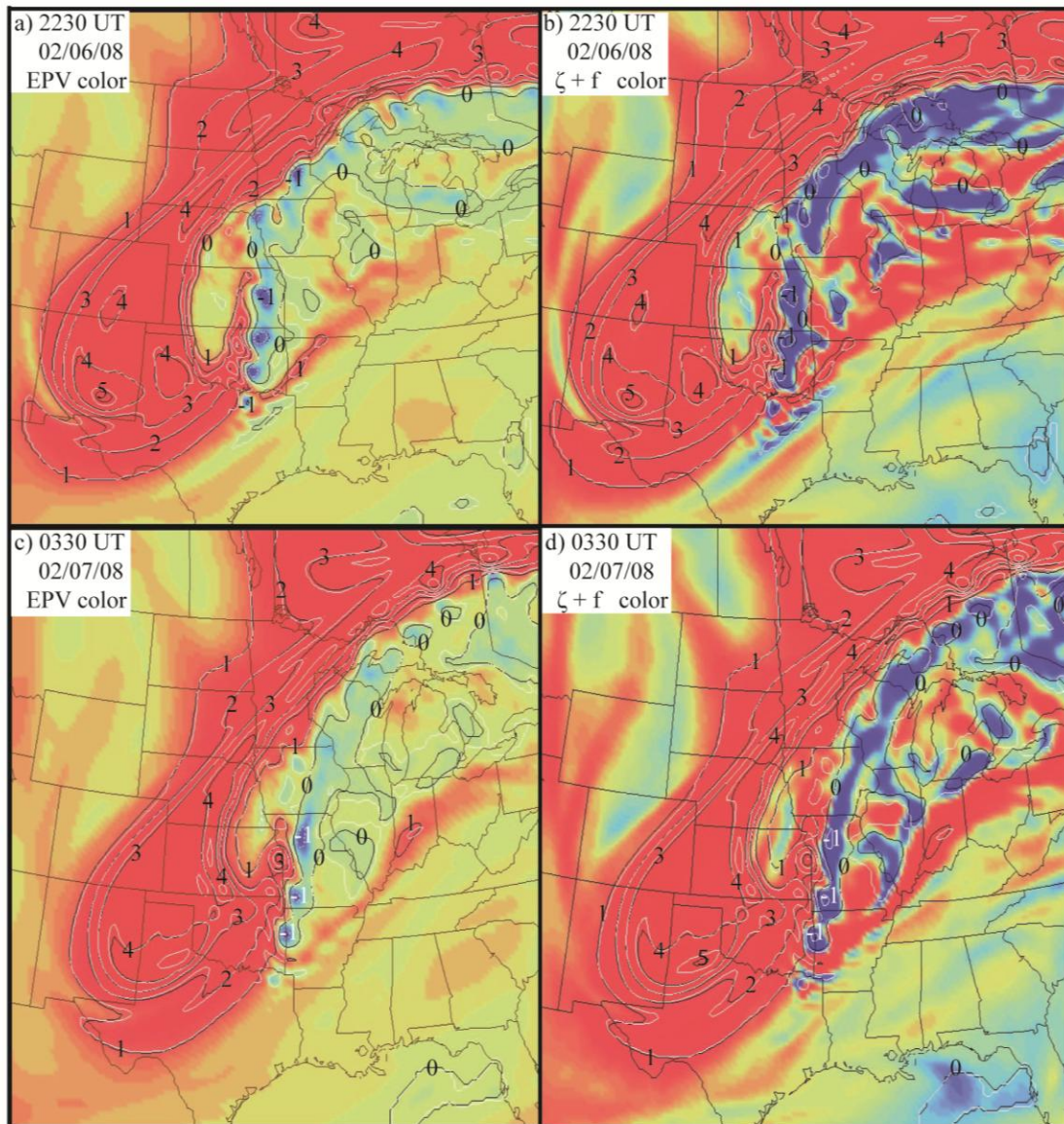


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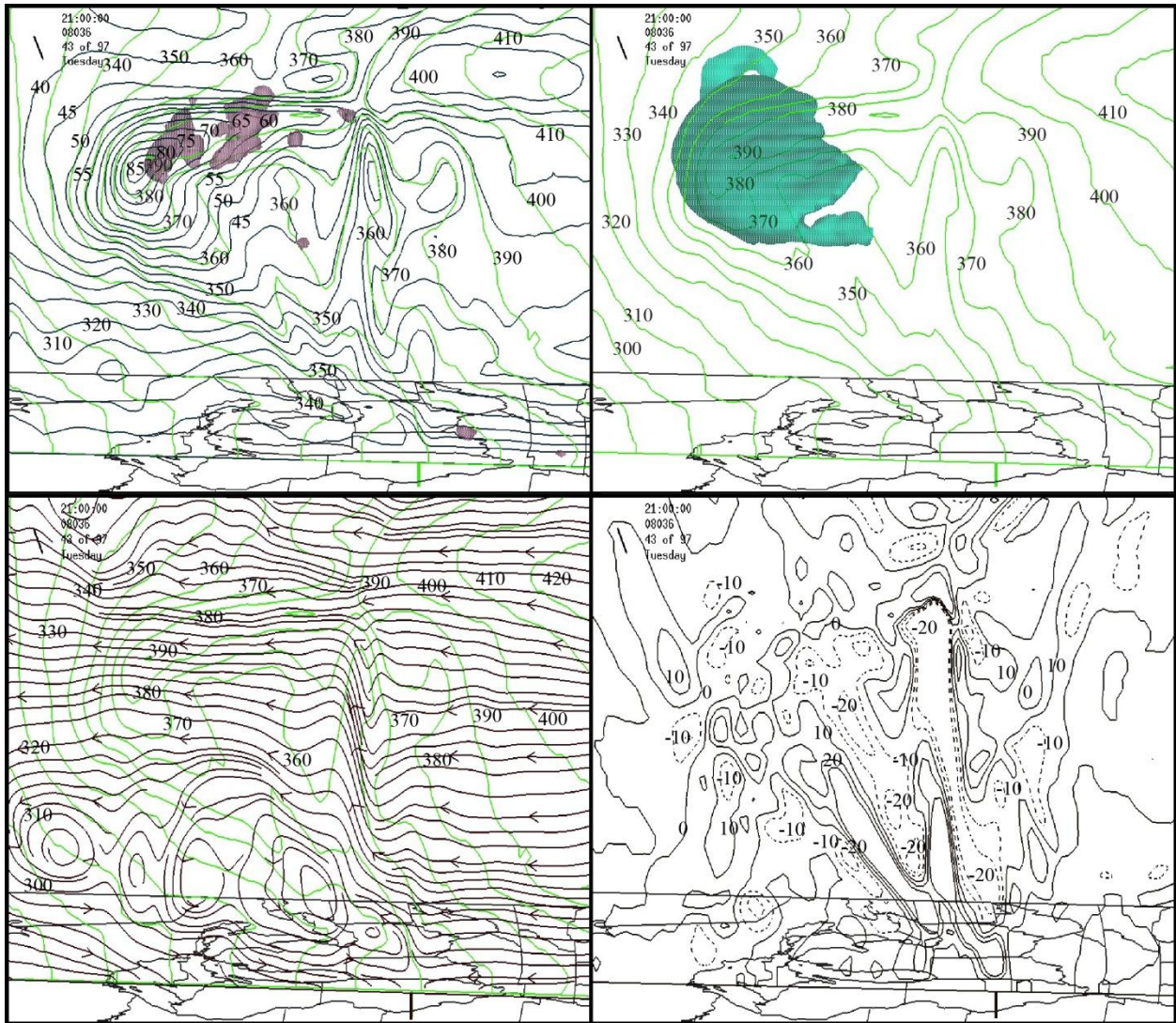
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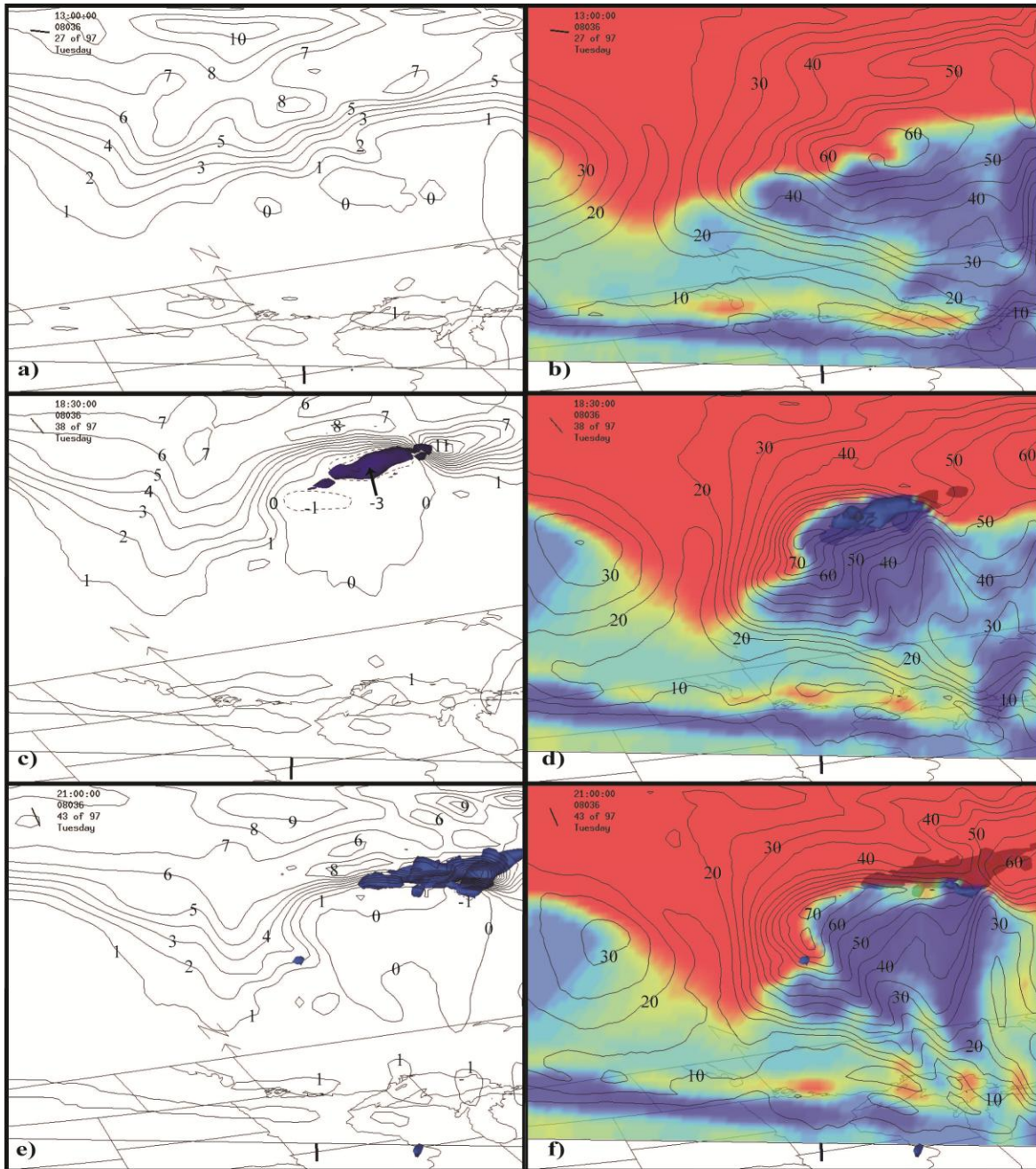


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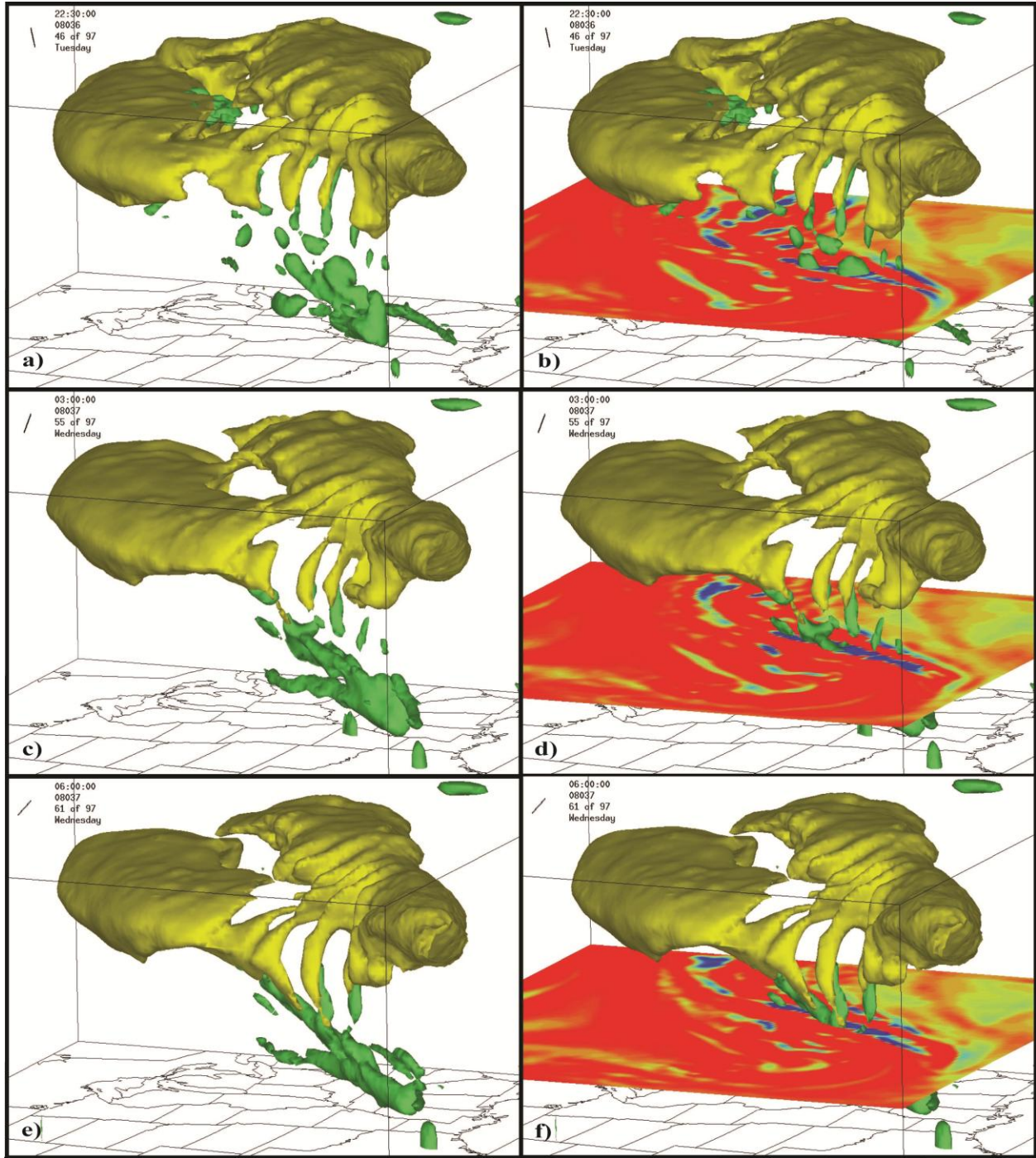
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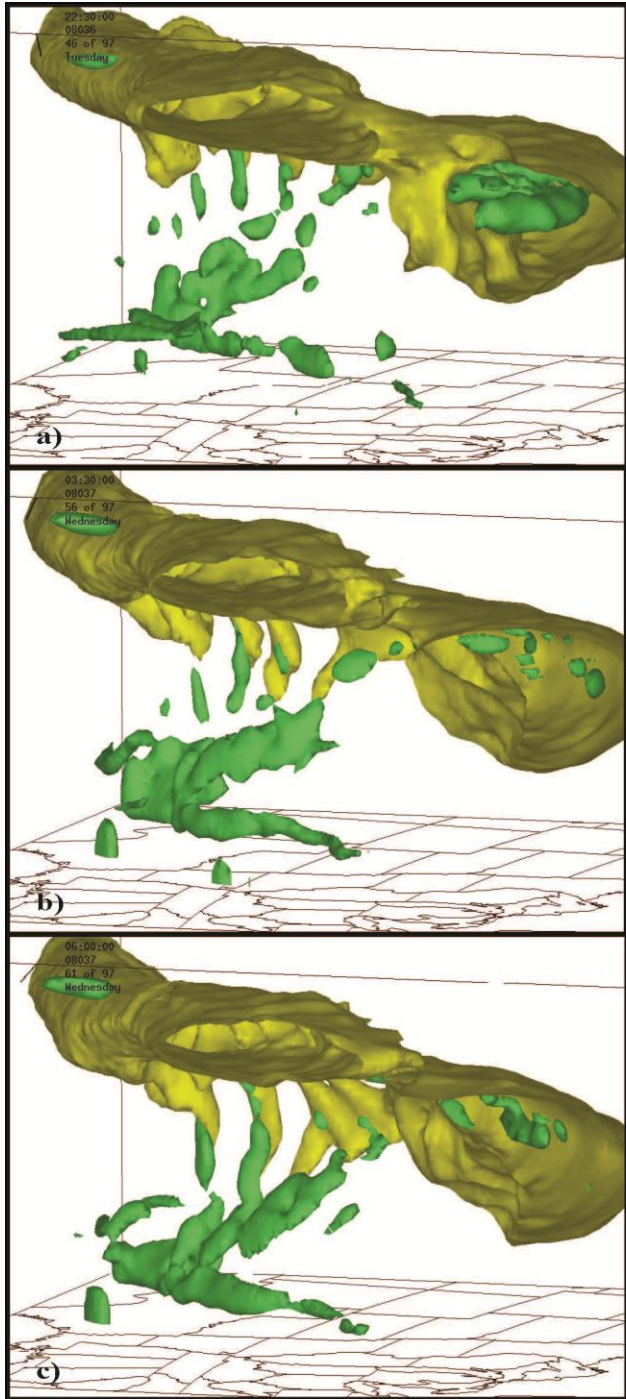
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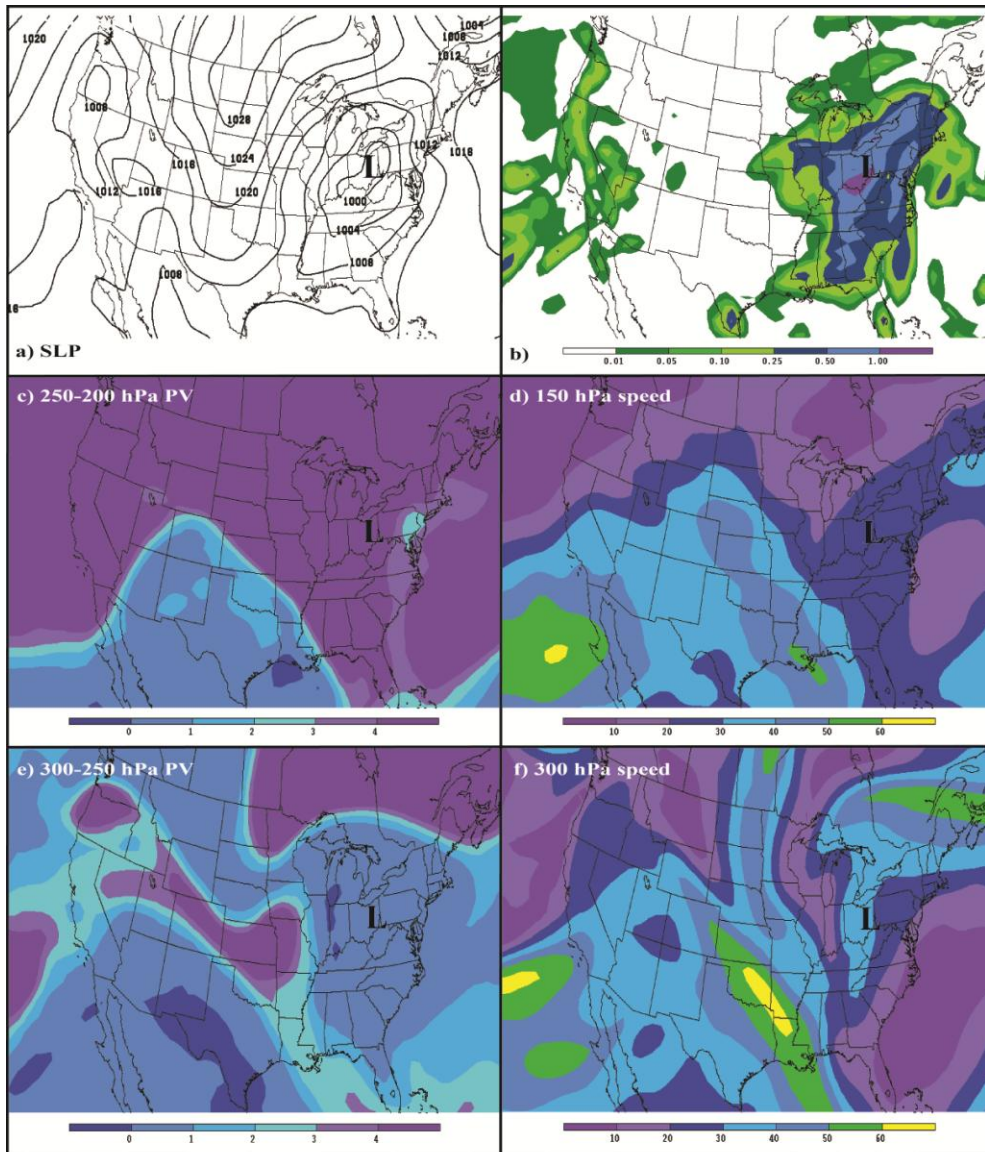
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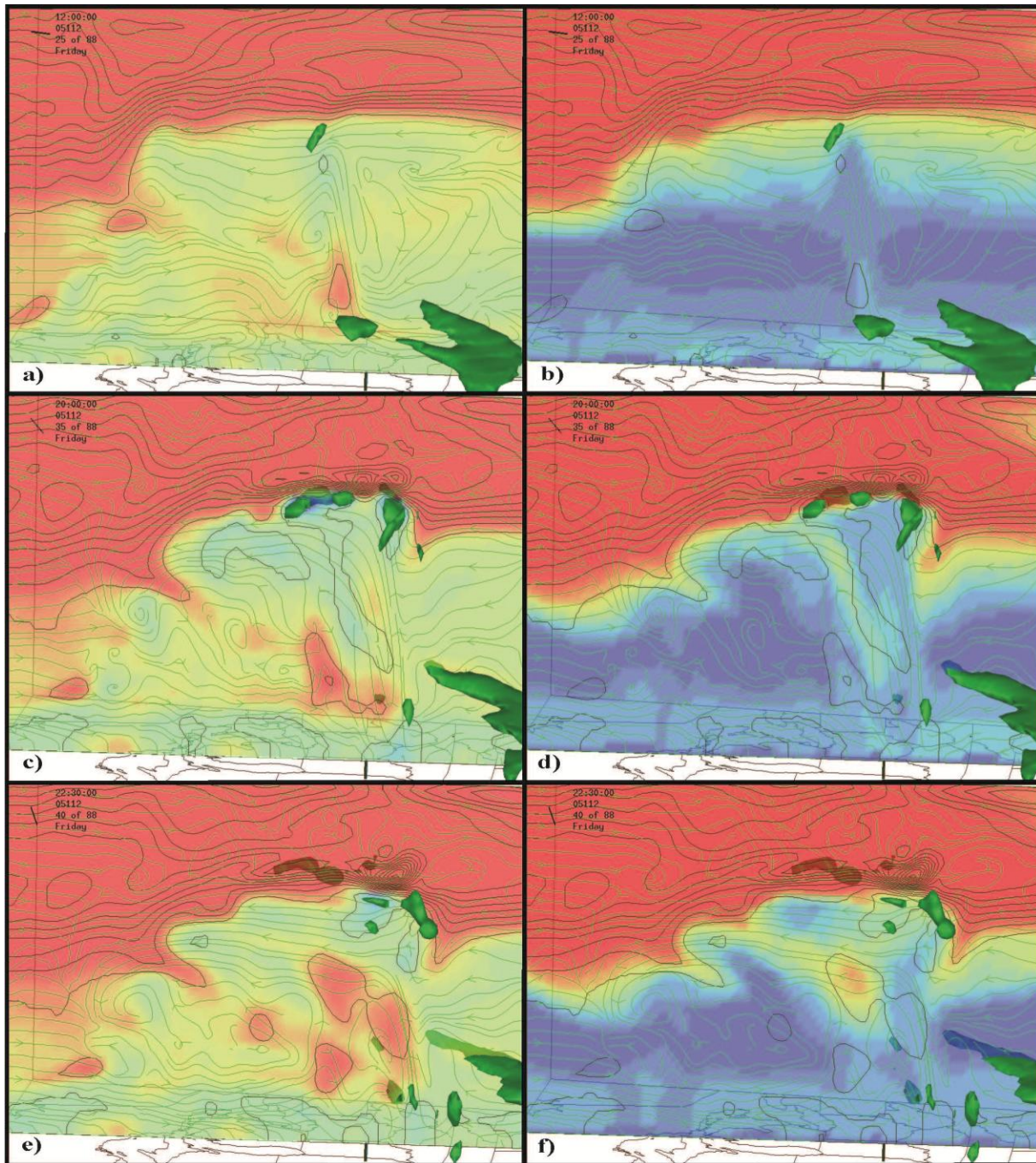
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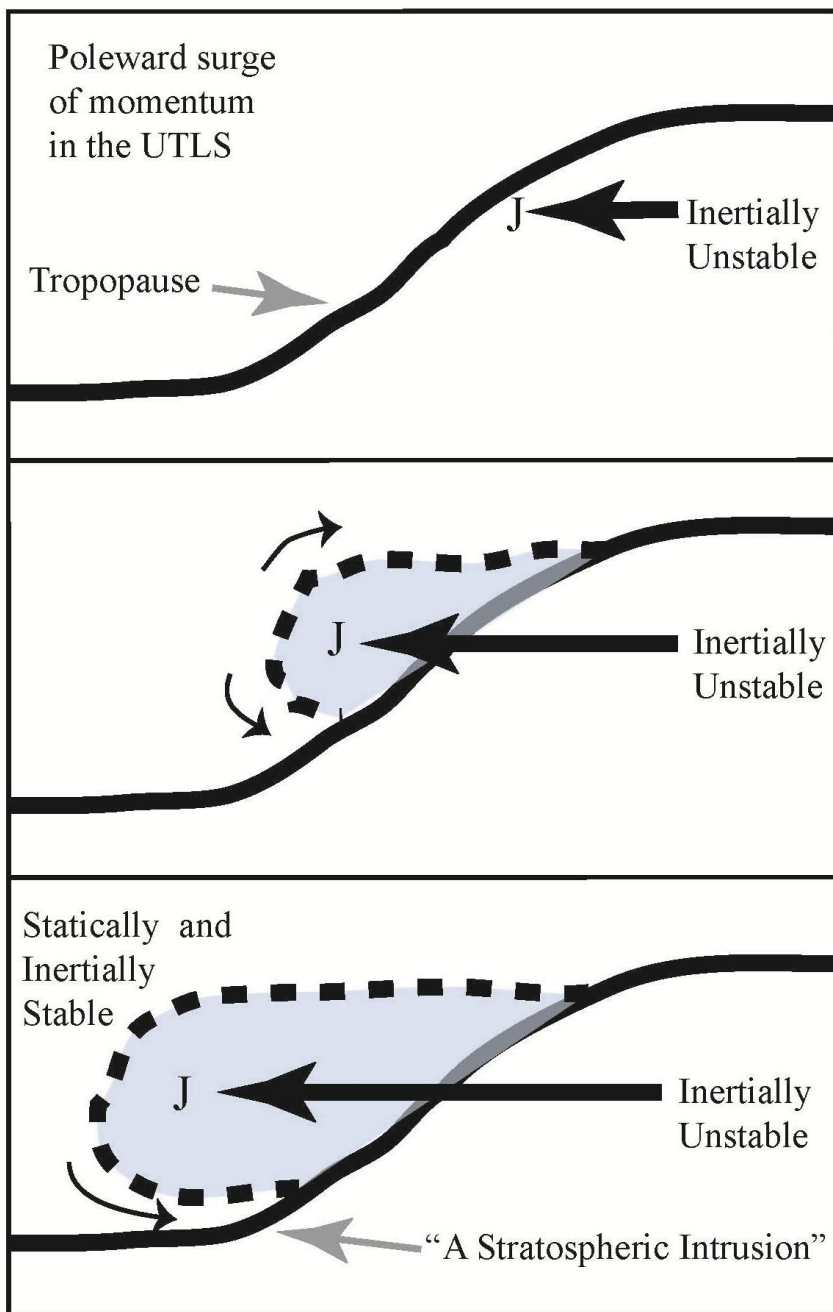
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